

Review exercise 2

1 Changing the units:

$$\text{diameter} = 7 \text{ cm} \Rightarrow \text{radius } (r) = 3.5 \text{ cm} = 0.035 \text{ m}$$

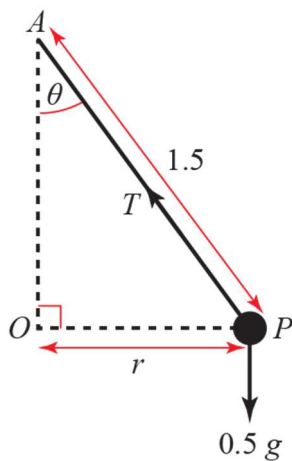
$$\text{angular speed } (\omega) = 1000 \text{ revolutions per minute} = \frac{1000 \times 2\pi}{60} \text{ radians per second}$$

Using $v = r\omega$ gives:

$$v = 0.035 \times \frac{1000 \times 2\pi}{60} = 3.67 \text{ ms}^{-1} \text{ (3 s.f.)}$$

So the speed is 3.67 ms^{-1} (3 s.f.)

2 a Let the tension in the string be T and let the string make an angle θ with the vertical.



Let $OP = r$, then $r = 1.5 \sin \theta$

Acceleration towards the centre of the circle $= r\omega^2 = 1.5 \sin \theta \times 2.7^2$

Force towards the centre of the circle $= T \sin \theta$

So using $F = ma$ gives:

$$T \sin \theta = 0.5 \times 1.5 \sin \theta \times 2.7^2$$

$$\Rightarrow T = 0.5 \times 1.5 \times 2.7^2 = 5.4675 = 5.5 \text{ N (2 s.f.)}$$

b Resolving the forces vertically $R(\uparrow)$ and substituting for T gives:

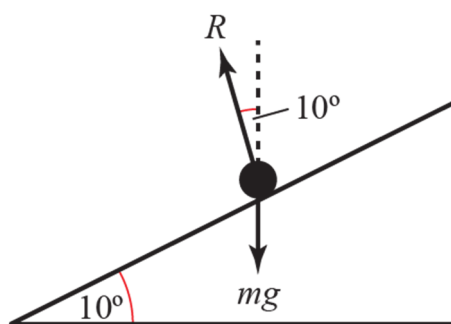
$$T \cos \theta - 0.5g = 0$$

$$\Rightarrow T \cos \theta = 0.5g$$

$$\Rightarrow \cos \theta = \frac{0.5g}{5.4675} = 0.8962$$

So $\theta = \cos^{-1} 0.8962 = 26^\circ$ (to the nearest degree)

3 The forces acting on the car are its weight mg and the normal reaction R .



$$R(\uparrow): R \cos 10^\circ - mg = 0$$

$$\Rightarrow R = \frac{mg}{\cos 10^\circ}$$

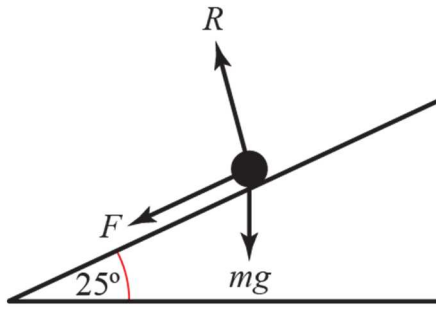
Using $F = ma$ horizontally gives:

$$R \sin 10^\circ = \frac{mv^2}{r} = \frac{m \times 18^2}{r}$$

$$\Rightarrow \frac{mg}{\cos 10^\circ} \times \sin 10^\circ = \frac{m \times 18^2}{r} \quad \text{substituting for } R$$

$$\Rightarrow r = \frac{18^2}{g \tan 10^\circ} = 187.4995\dots = 190 \text{ m (2 s.f.)}$$

- 4 The forces acting on the cyclist and bicycle are their weight mg , the normal reaction R and the friction acting down the slope.



$$R(\uparrow): R \cos 25^\circ - F \sin 25^\circ - mg = 0$$

If μ is the coefficient of friction between the cycle's tyres and the track, then the maximum friction for which the tyres do not slip is $F = \mu R = 0.6R$. Substituting for F gives:

$$R \cos 25^\circ - 0.6R \sin 25^\circ - mg = 0$$

$$\Rightarrow R = \frac{mg}{\cos 25^\circ - 0.6 \sin 25^\circ} \quad (1)$$

$$R(\leftarrow): R \sin 25^\circ + F \cos 25^\circ = \frac{mv^2}{r} \quad \text{using } F = ma \text{ and } a = \frac{v^2}{r}$$

$$\Rightarrow R \sin 25^\circ + 0.6R \cos 25^\circ = \frac{mv^2}{40} \quad \text{as } r = 40 \text{ and } F = 0.6R \text{ at maximum speed}$$

$$\Rightarrow R = \frac{mv^2}{40(\sin 25^\circ + 0.6 \cos 25^\circ)} \quad (2)$$

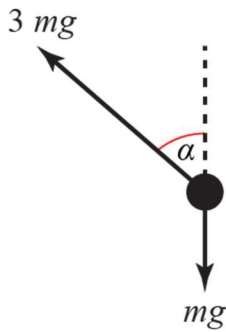
Using equations (1) and (2) gives:

$$\frac{mv^2}{40(\sin 25^\circ + 0.6 \cos 25^\circ)} = \frac{mg}{\cos 25^\circ - 0.6 \sin 25^\circ}$$

$$\Rightarrow v^2 = \frac{40g(\sin 25^\circ + 0.6 \cos 25^\circ)}{\cos 25^\circ - 0.6 \sin 25^\circ} = 580.37 \text{ (2 d.p.)}$$

$$\Rightarrow v = 24 \text{ m s}^{-1} \text{ (2 s.f.)}$$

- 5 a The forces acting on the metal ball are the weight of the ball and tension along the string.



$$R(\uparrow): 3mg \cos \alpha - mg = 0$$

$$\Rightarrow \cos \alpha = \frac{mg}{3mg} = \frac{1}{3}$$

$$\Rightarrow \alpha = 70.5^\circ \text{ (3 s.f.)}$$

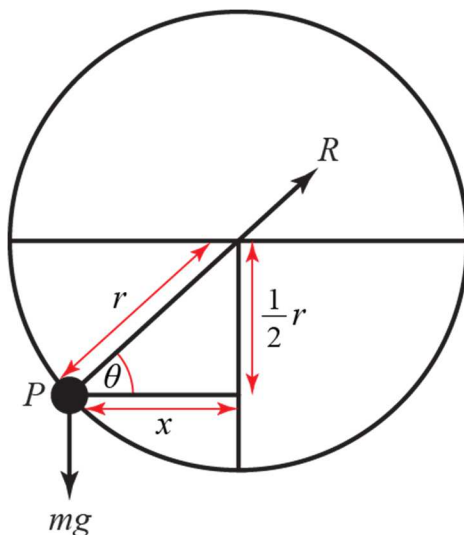
- b $R(\leftarrow): 3mg \sin \alpha = mr2gk$ using $F = ma$ and $a = r\omega^2$

Let the length of the string be l , then from $\triangle AOB$ it is clear that $\sin \alpha = \frac{r}{l}$

$$\text{So } 3mg \frac{r}{l} = mr2gk$$

$$\Rightarrow l = \frac{3}{2k}$$

- 6 a The forces acting on the particle are its weight and the normal reaction.



Let θ be the angle between the normal reaction and the horizontal.

$$\text{Then } \sin \theta = \frac{\frac{1}{2}r}{r} = \frac{1}{2} \Rightarrow \theta = 30^\circ$$

$$R(\uparrow): R \sin 30^\circ - mg = 0$$

$$\Rightarrow R = \frac{mg}{\sin 30^\circ} = 2mg$$

6 b $R(\rightarrow): R \cos 30^\circ = m x \omega^2$ using $F = ma$ and $a = r\omega^2$

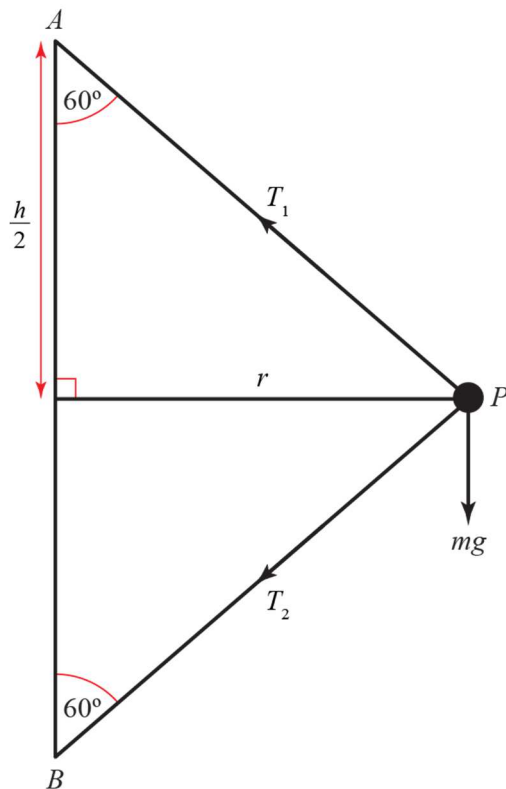
Using the result from part a and as $x = r \cos 30^\circ$ this gives:

$$2mg = mr\omega^2$$

$$\Rightarrow \omega = \sqrt{\frac{2g}{r}}$$

$$\text{Time to complete one revolution} = \frac{2\pi}{\omega} = 2\pi \sqrt{\frac{r}{2g}}$$

- 7 a Let T_1 be the tension in AP and T_2 be the tension in BP . The forces acting on P are the tensions in the two strings and its weight.



From the diagram, it can be seen that the equilateral triangle APB can be divided into two right-angled triangles, where:

$$\tan 60^\circ = \frac{r}{\frac{h}{2}}$$

$$\Rightarrow r = \frac{h}{2} \times \tan 60^\circ = \frac{\sqrt{3}h}{2}$$

7 b Resolving the forces from the diagram in part a

$$R(\uparrow): T_1 \cos 60^\circ - T_2 \cos 60^\circ - mg = 0$$

$$\Rightarrow T_1 - T_2 = 2mg \quad (1)$$

$$R(\leftarrow): T_1 \sin 60^\circ + T_2 \sin 60^\circ = mr\omega^2 \quad \text{using } F = ma \text{ and } a = r\omega^2$$

$$\Rightarrow \frac{\sqrt{3}}{2}T_1 + \frac{\sqrt{3}}{2}T_2 = m \frac{\sqrt{3}}{2}h\omega^2 \quad \text{as } r = \frac{\sqrt{3}}{2}h \text{ from part a}$$

$$\Rightarrow T_1 + T_2 = mh\omega^2 \quad (2)$$

Adding equations (1) and (2) gives: $2T_1 = 2mg + mh\omega^2 \Rightarrow T_1 = mg + \frac{1}{2}mh\omega^2$

Substituting for T_1 in equation (2) gives: $T_2 = \frac{1}{2}mh\omega^2 - mg$

c Both strings are taut, therefore $T_1 > 0$ and $T_2 > 0$. From part b, $T_1 > 0$ for all values of ω

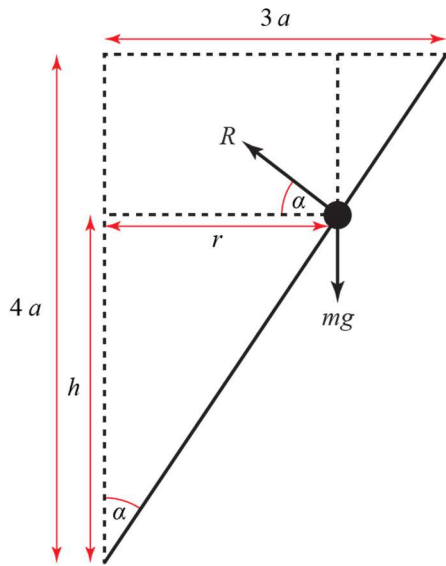
From the formula for T_2 in part b, the condition that $T_2 > 0 \Rightarrow \omega > \sqrt{\frac{2g}{h}}$

$$\text{As } T = \frac{2\pi}{\omega}, \omega = \frac{2\pi}{T}$$

$$\text{So } \omega > \sqrt{\frac{2g}{h}} \Rightarrow \frac{2\pi}{T} > \sqrt{\frac{2g}{h}} \Rightarrow T < 2\pi\sqrt{\frac{h}{2g}} \quad \text{as } T > 0 \text{ and } \sqrt{\frac{h}{2g}} < 0$$

$$\Rightarrow T < \pi\sqrt{\frac{2h}{g}} \quad \text{as required}$$

- 8 Let α be the angle between the slant side and the axis of the cone, r the radius of the horizontal circle with centre C and h the height of C above V . The forces acting on the particle are its weight and the normal reaction.



From the diagram, $\tan \alpha = \frac{3a}{4a} = \frac{3}{4}$ and $\tan \alpha = \frac{r}{h}$

$$R(\uparrow): R \sin \alpha - mg = 0$$

$$\Rightarrow R \sin \alpha = mg \quad (1)$$

$$R(\leftarrow): R \cos \alpha = mr\omega^2 \quad \text{using } F = ma \text{ and } a = r\omega^2$$

$$\Rightarrow R \cos \alpha = \frac{mr8g}{9a} \quad (2) \quad \text{using } \omega = \sqrt{\frac{8g}{9a}}$$

Dividing equation (1) by equation (2) gives:

$$\tan \alpha = mg \div \frac{8mrg}{9a} = \frac{9a}{8r}$$

$$\Rightarrow \frac{3}{4} = \frac{9a}{8r} \quad \text{using } \tan \alpha = \frac{3}{4}$$

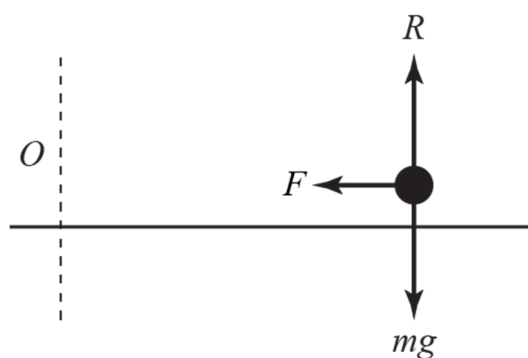
$$\Rightarrow r = \frac{3a}{2}$$

$$\text{As } \tan \alpha = \frac{3}{4} = \frac{r}{h}, h = \frac{4}{3}r$$

$$\text{So } h = \frac{4}{3} \times \frac{3a}{2} = 2a$$

The height of C above V is $2a$.

- 9 a The forces acting on the particle are its weight, the normal reaction and friction.



$$R(\uparrow): R - mg = 0 \Rightarrow R = mg$$

$$R(\leftarrow): F = mr\omega^2 \quad \text{using } F = ma \text{ and } a = r\omega^2$$

$$\Rightarrow F = \frac{4ma\omega^2}{3} \quad \text{as } r = \frac{4}{3}a$$

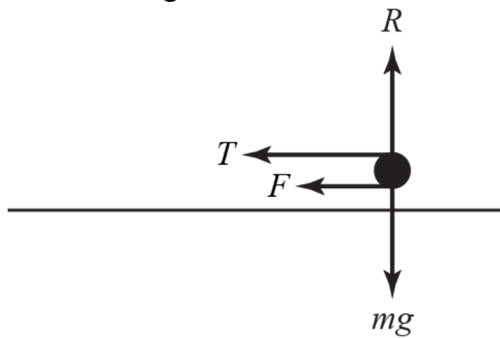
As P remains at rest $F \leq \mu R$

$$\text{So } \frac{4ma\omega^2}{3} \leq \frac{3}{5}R$$

$$\Rightarrow \frac{4ma\omega^2}{3} \leq \frac{3mg}{5}$$

$$\Rightarrow \omega^2 \leq \frac{9g}{20a} \quad \text{as required}$$

- 9 b The forces now acting on the particle are its weight, the normal reaction, friction and the tension in the elastic string.



Find T using Hooke's law (Further Mechanics 1, Chapter 3): $T = \frac{\lambda x}{l}$,

where λ is the modulus of elasticity, x is the extension of the string and l is its natural length

So in this case, $T = \frac{2mg}{a} \times \frac{a}{3} = \frac{2mg}{3}$

$R(\leftarrow): T + F = mr\omega^2$ using $F = ma$ and $a = r\omega^2$

$$\Rightarrow \omega^2 = \frac{3}{4ma} \left(\frac{2mg}{3} + F \right)$$

Note that the frictional force can act away from O against the pull of the elastic string, or towards O through the force of the acceleration of the particle generated by the circular motion. As P remains at rest $-\mu R \leq F \leq \mu R$ (where $\mu R = 0.6mg$ from part a).

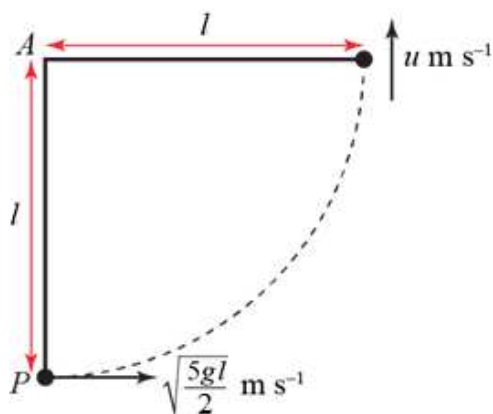
So, from the equation for ω^2 , it is maximum when $F = \mu R = \frac{3mg}{5}$

$$\text{So } \omega_{\max}^2 = \frac{3}{4ma} \left(\frac{2mg}{3} + \frac{3mg}{5} \right) = \frac{3}{4ma} \times \frac{19mg}{15} = \frac{19g}{20a}$$

Similarly ω^2 is minimum when $F = -\mu R = -\frac{3mg}{5}$

$$\text{So } \omega_{\min}^2 = \frac{3}{4ma} \left(\frac{2mg}{3} - \frac{3mg}{5} \right) = \frac{3}{4ma} \times \frac{mg}{15} = \frac{g}{20a}$$

10 a Let u be the speed of P when the string is horizontal.



Applying the work-energy principle, the sum of the particle's kinetic and gravitational potential energy remains constant, so when it reaches the horizontal the loss of kinetic energy = the gain in potential energy.

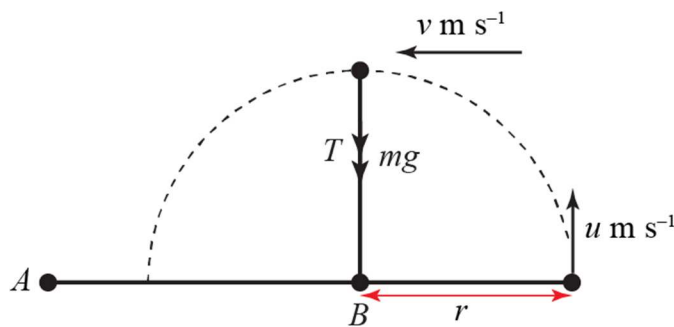
Take the starting point as the zero level for potential energy.

$$\text{So } \frac{1}{2}m \frac{5gl}{2} - \frac{1}{2}mu^2 = mgl$$

$$\Rightarrow u^2 = \frac{5gl}{2} - 2gl = \frac{gl}{2}$$

$$\Rightarrow u = \sqrt{\frac{gl}{2}}$$

10 b Let the particle move in a semi-circle about B with radius r .



Taking the line AB as the zero level for potential energy, then applying conservation of energy at the highest point of the semi-circle:

$$\frac{1}{2}m(u^2 - v^2) = mgr$$

$$\Rightarrow v^2 = u^2 - 2gr \quad (1)$$

Resolving the vertical forces at the highest point:

$$R(\downarrow): T + mg = \frac{mv^2}{r} \quad \text{using } F = ma \text{ and } a = \frac{v^2}{r}$$

$$\Rightarrow T = \frac{m(u^2 - 2gr)}{r} - mg \quad \text{substituting for } v^2 \text{ using equation (1)}$$

$$\Rightarrow T = \frac{mu^2}{r} - 3mg$$

$$\Rightarrow T = \frac{mgl}{2r} - 3mg \quad \text{s substituting for } u^2 \text{ using result from part a}$$

As the string does not go slack $T > 0$, so

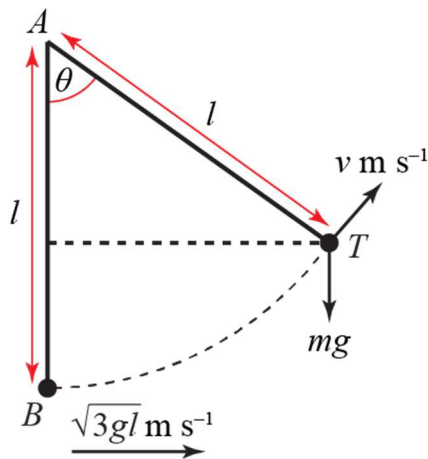
$$\frac{mgl}{2r} - 3mg > 0$$

$$\Rightarrow mgl > 6mgr$$

$$\Rightarrow r < \frac{l}{6}$$

As $AB = l - r$, this shows that $AB > \frac{5l}{6}$

11 a Let the speed of the particle be v when it makes an angle θ with the downward vertical.



Take the horizontal through B as the zero level for potential energy. Using conservation of energy:

$$\frac{1}{2}m(3gl - v^2) = mgl(1 - \cos \theta)$$

$$\Rightarrow v^2 = 3gl - 2gl(1 - \cos \theta)$$

$$\Rightarrow v^2 = gl + 2gl \cos \theta \quad (1)$$

Resolving along the string:

$$R(\curvearrowright): T - mg \cos \theta = \frac{mv^2}{l}$$

$$\text{using } F = ma \text{ and } a = \frac{v^2}{r} \text{ and } r = l$$

$$\Rightarrow T = mg \cos \theta + \frac{mgl + 2gl \cos \theta}{l}$$

substituting for v^2 using equation (1)

$$\Rightarrow T = mg + 3mg \cos \theta = mg(1 + 3 \cos \theta)$$

as required

b The instant the string becomes slack, $T = 0$.

Using the expression for T from part **a**, this occurs when:

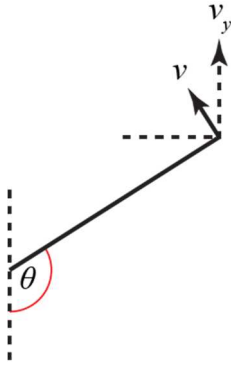
$$1 + 3 \cos \theta = 0 \Rightarrow \cos \theta = -\frac{1}{3}$$

Substituting for $\cos \theta$ in equation (1) gives:

$$v^2 = gl + 2gl \times -\frac{1}{3} = \frac{gl}{3}$$

$$\Rightarrow v = \sqrt{\frac{gl}{3}}$$

- 11 c** The particle now moves as a projectile under gravity. The maximum height is achieved when the vertical component of the velocity is zero.



$$v_y = v \sin(180^\circ - \theta) = \sqrt{\frac{gl}{3}} \sin \theta$$

$$\text{If } \frac{\pi}{2} < \theta < \pi \text{ and } \cos \theta = -\frac{1}{3}, \text{ then } \sin \theta = \frac{2\sqrt{2}}{3}$$

$$\text{So } v_y = \frac{2\sqrt{2}}{3} \sqrt{\frac{gl}{3}}$$

Considering the particle's vertical motion, $u = v_y$, $v = 0$, $a = -g$, $s = h$, where h is the height above the point the string becomes slack.

$$\text{Using } v^2 = u^2 - 2gh$$

$$h = \frac{v_y^2}{2g} = \frac{8}{9} \times \frac{gl}{3} \times \frac{1}{2g} = \frac{4l}{27}$$

$$\text{The height at which string becomes slack} = l + l \cos(180^\circ - \theta) = l(1 - \cos \theta) = \frac{4l}{3}$$

So if H is the maximum height above the level of B reached by P

$$H = \frac{4l}{3} + \frac{4l}{27} = \frac{40l}{27}$$

- 12 a** Take the horizontal through l as the zero level for potential energy. When angle $\theta = \alpha$, the particle is momentarily at rest. Using conservation of energy, with the loss of kinetic energy equal to the gain in potential energy:

$$\frac{1}{2} mu^2 = mgl(1 - \cos \alpha) = \frac{mgl}{3} \quad \text{as } \cos \alpha = \frac{2}{3}$$

$$\Rightarrow u^2 = \frac{2}{3} gl$$

$$\Rightarrow u = \sqrt{\frac{2gl}{3}}$$

- 12 b** Let the speed of the particle at θ be v .
Resolving along the string gives:

$$R(\nearrow): T - mg \cos \theta = \frac{mv^2}{l}$$

$$\Rightarrow T = mg \cos \theta + \frac{mv^2}{l}$$

$$\text{using } F = ma \text{ and } a = \frac{v^2}{r} \text{ and } r = l$$

Applying conservation of energy:

$$\frac{1}{2}mu^2 - \frac{1}{2}mv^2 = mgl(1 - \cos \theta)$$

$$\Rightarrow v^2 = u^2 - 2gl(1 - \cos \theta) = \frac{2gl}{3} - 2gl(1 - \cos \theta)$$

using the result from part a

$$\Rightarrow v^2 = 2gl \cos \theta - \frac{4gl}{3}$$

Substituting for v^2 in the equation for T gives:

$$T = mg \cos \theta + 2mg \cos \theta - \frac{4mg}{3}$$

$$= 3mg \cos \theta - \frac{4mg}{3}$$

$$= \frac{mg}{3}(9 \cos \theta - 4)$$

- c** Maximum value of T is when $\cos \theta = 1$

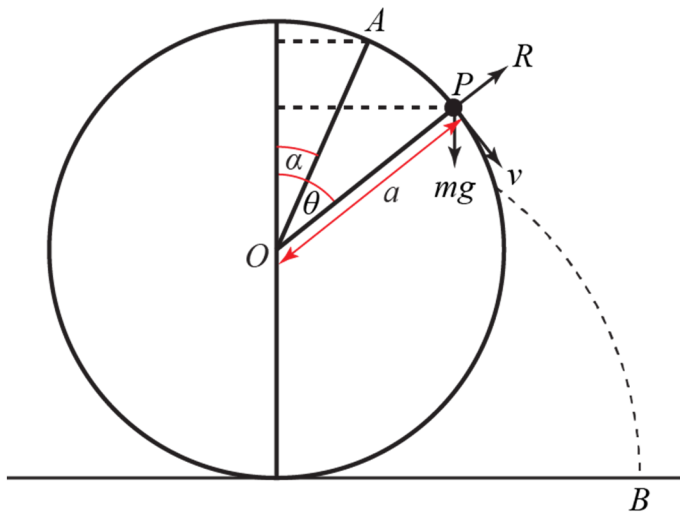
$$\text{So } T_{\max} = \frac{5mg}{3}$$

Minimum value of T is when $\cos \theta = \frac{2}{3}$

$$\text{So } T_{\min} = \frac{2mg}{3}$$

$$\text{Hence } \frac{2mg}{3} \leq T \leq \frac{5mg}{3}$$

13 a This is a diagram of the problem.



Take the horizontal through A as the zero level for potential energy. Using conservation of energy, the gain of kinetic energy at P equals the loss in potential energy:

$$\frac{1}{2}mv^2 = mg(a \cos \alpha - a \cos \theta)$$

$$v^2 = 2ga(\cos \alpha - \cos \theta)$$

b Resolving along the radius OP :

$$R (\sphericalangle): mg \cos \theta - R = \frac{mv^2}{a}$$

At the point when P loses contact with the sphere $R = 0$

$$\Rightarrow mg \cos \theta = \frac{mv^2}{a} = 2gm(\cos \alpha - \cos \theta)$$

substituting for v^2 from part a

$$\Rightarrow 3 \cos \theta = 2 \cos \alpha = \frac{3}{2}$$

$$\Rightarrow \cos \theta = \frac{1}{2}, \text{ so } \theta = 60^\circ \text{ or } \frac{\pi}{3} \text{ radians}$$

c Let the speed of P when it hits the table be w . Then taking the horizontal through A as the zero level for potential energy and using conservation of energy, the gain of kinetic energy at B equals the loss in potential energy:

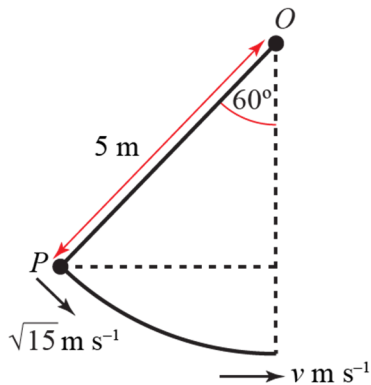
$$\frac{1}{2}mw^2 = mg(a + a \cos \alpha)$$

$$\Rightarrow w^2 = 2ga \left(1 + \frac{3}{4} \right) = \frac{7ga}{2}$$

$$\text{So } w = \sqrt{\frac{7ga}{2}} \text{ ms}^{-1}$$

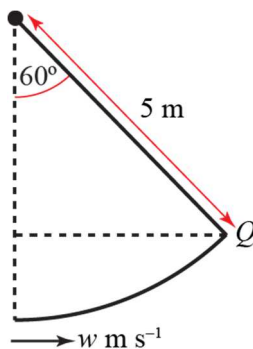
There are alternative ways to calculate w that use projectiles, but the approach shown above is the shortest method.

- 14 a** Let the speed of the trapeze artist at the lowest point of her path be v . Then taking the horizontal through A as the zero level for potential energy and using conservation of energy, the gain in kinetic energy at this point equals the loss in potential energy:



$$\begin{aligned}\frac{1}{2}mv^2 - \frac{1}{2}m(\sqrt{15})^2 &= mg \times 5(1 - \cos 60^\circ) \\ \Rightarrow v^2 &= 15 + 5g = 64 \\ \Rightarrow v &= 8 \text{ m s}^{-1}\end{aligned}$$

- b** Let the velocity after catching the ball be w . Then the loss in kinetic energy at point Q where the trapeze artist becomes momentarily stationary equals the gain in potential energy from the lowest point of the trapeze artist's path.



$$\begin{aligned}\frac{1}{2}(60 + m)w^2 - 0 &= (60 + m)g5(1 - \cos 60^\circ) \\ \Rightarrow w^2 &= 5g = 49, \text{ so } w = 7 \text{ m s}^{-1}\end{aligned}$$

So at the instant just prior to catching the ball, the trapeze artist is travelling at 8 m s^{-1} (part **a**) and the ball is travelling at 3 m s^{-1} in the opposite direction, and immediately after catching the ball she is travelling at 7 m s^{-1} . Using conservation of linear momentum at the instant when she catches the ball, this gives:

$$\begin{aligned}60 \times 8 + (-3 \times m) &= (60 + m) \times 7 \\ \Rightarrow 480 - 3m &= 420 + 7m \\ \Rightarrow 10m &= 60 \\ \text{So } m &= 6 \text{ kg}\end{aligned}$$

- c** $R(\uparrow): T - 66g = \frac{66w^2}{r}$ using $F = ma$ and $a = \frac{v^2}{r}$
- $$\Rightarrow T = 66g + 66 \times \frac{7^2}{5} = 1293.6 = 1300 \text{ N (2 s.f.)}$$

- 15 a** Let the speed at C be $v \text{ ms}^{-1}$. Then taking the horizontal through B as the zero level for potential energy and using conservation of energy, the gain in kinetic energy at point C equals the loss in potential energy:

$$\frac{1}{2}mv^2 - \frac{1}{2}m \times 20^2 = mg \times 50(1 - \cos 60^\circ)$$

$$\Rightarrow v^2 = 20^2 + 50g \Rightarrow v^2 = 890$$

$$\text{So } v = 30 \text{ ms}^{-1} \text{ (2 s.f.)}$$

- b** At C , $R(\uparrow)$: $R - 70g = \frac{70v^2}{50}$ using $F = ma$ and $a = \frac{v^2}{r}$

$$\Rightarrow R = 70g + \frac{70 \times 890}{50} = 686 + 1246 = 1900 \text{ N (2 s.f.)}$$

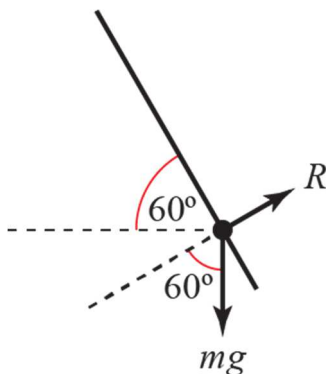
- c** Consider motion C to D . Let the speed at D be $w \text{ ms}^{-1}$. Then taking the horizontal through C as the zero level for potential energy and using conservation of energy, the loss in kinetic energy at point C equals the gain in potential energy:

$$\frac{1}{2}m \times 890 - \frac{1}{2}mw^2 = mg \times 50(1 - \cos 30^\circ)$$

$$\Rightarrow w^2 = 890 - 100g(1 - \cos 30^\circ) = 759 \text{ (3 s.f.)}$$

$$\Rightarrow w = 28 \text{ ms}^{-1} \text{ (2 s.f.)}$$

- d** Resolving perpendicular to the slope at B just before the circular motion and just after circular motion begins:



$$\text{Before: } R = mg \cos 60^\circ = 35g$$

$$\text{After: } R - mg \cos 60^\circ = \frac{m \times 20^2}{50} \Rightarrow R = 35g + 560$$

$$\text{So change in } R = 560 \text{ N}$$

- e** Allowing for the influence of friction would mean that the skier would arrive at C with lower speed. From the equation used in part **b**, this would result in a lower normal reaction.

- 16 a** Let the velocity of the particle at C be v . Taking the horizontal through A as the zero level for potential energy and using conservation of energy, the gain in kinetic energy at C equals the loss in potential energy:

$$\begin{aligned}\frac{1}{2}mv^2 - \frac{1}{2}mu^2 &= mga(1 - \cos\theta) \\ \Rightarrow v^2 &= u^2 + 2ga(1 - \cos\theta) \quad (1)\end{aligned}$$

Resolving along the radius through C :

$$R (\sphericalangle): -R + mg \cos\theta = \frac{mv^2}{a}$$

As the particle leaves the sphere at C , $R = 0$ so $v^2 = ag \cos\theta$

Substituting for v^2 in equation (1) gives:

$$ag \cos\theta = u^2 + 2ga(1 - \cos\theta)$$

$$\Rightarrow 3 \cos\theta = \frac{u^2}{ag} + 2$$

$$\Rightarrow \cos\theta = \frac{2}{3} + \frac{u^2}{3ag}$$

- b** Using conservation of energy, the gain in kinetic energy from C to when the particle hits the ground equals the loss in potential energy:

$$\frac{1}{2}m\left(\frac{9ag}{2}\right) - \frac{1}{2}m(ag \cos\theta) = mga(1 + \cos\theta)$$

$$\Rightarrow \frac{3}{2}\cos\theta = \frac{9}{4} - 1 = \frac{5}{4} \Rightarrow \cos\theta = \frac{5}{6}$$

So $\theta = 34^\circ$ (2 s.f.)

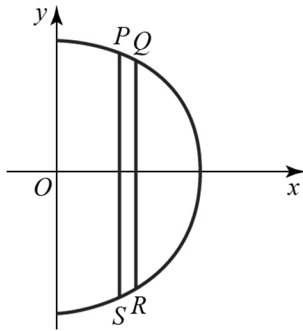
- 17** The centre of mass lies on the axis of symmetry, $y = 0$.

To find the x coordinate, using $\bar{x} = \frac{\int y^2 x \, dx}{\int y^2 \, dx}$, and substituting $y = \sqrt{x}$ gives:

$$\bar{x} = \frac{\int_0^4 x^2 \, dx}{\int_0^4 x \, dx} = \frac{\left[\frac{1}{3}x^3\right]_0^4}{\left[\frac{1}{2}x^2\right]_0^4} = \frac{64}{3} \div 8 = \frac{8}{3}$$

So the coordinates of the centre of mass of the solid are $\left(\frac{8}{3}, 0\right)$, a distance of $\frac{8}{3}$ from O .

- 18 a** Let the straight edge lie along the y -axis. Then the centre of mass lies on the x -axis from symmetry.



If P has coordinates (x, y) and the elemental strip $PQRS$ has width δx then its area is $2y\delta x$.
The mass M of the lamina = $\frac{1}{2}\pi a^2\rho$, where ρ is the mass per unit area of the lamina.

Let \bar{x} be the distance of the centre of mass from O , then $M\bar{x} = \int_0^a 2\rho xy dx$

As the boundary of the semicircle has the equation $x^2 + y^2 = a^2$, then $y = (a^2 - x^2)^{\frac{1}{2}}$

$$\text{So } M\bar{x} = \int_0^a 2\rho x(a^2 - x^2)^{\frac{1}{2}} dx = \left[-\frac{2}{3}\rho(a^2 - x^2)^{\frac{3}{2}} \right]_0^a = \frac{2}{3}\rho a^3$$

$$\Rightarrow \bar{x} = \frac{\frac{2}{3}\rho a^3}{\frac{1}{2}\pi a^2\rho} = \frac{4a}{3\pi} \quad \text{as required}$$

- b** Using the result from part **a** and letting \bar{x} be the distance of the centre of mass of the resulting lamina from O :

Shape	Mass	Distance of centre of mass from O
Semicircle radius a	$\frac{1}{2}\pi\rho a^2$	$\frac{4a}{3\pi}$
Semicircle radius b	$\frac{1}{2}\pi\rho b^2$	$\frac{4b}{3\pi}$
Resulting	$\frac{1}{2}\pi\rho(a^2 - b^2)$	\bar{x}

Taking moments about O :

$$\frac{1}{2}\pi\rho(a^2 - b^2)\bar{x} = \frac{1}{2}\pi\rho a^2 \times \frac{4a}{3\pi} - \frac{1}{2}\pi\rho b^2 \times \frac{4b}{3\pi}$$

$$\Rightarrow \bar{x} = \frac{4}{3\pi} \frac{(a^3 - b^3)}{(a^2 - b^2)} = \frac{4}{3\pi} \frac{(a-b)(a^2 + ab + b^2)}{(a-b)(a+b)} = \frac{4}{3\pi} \frac{(a^2 + ab + b^2)}{(a+b)} \quad \text{as required}$$

- c** As $b \rightarrow a$, the area becomes a circular arc and from the equation found in part **b**

$$\bar{x} \rightarrow \frac{4}{3\pi} \times \frac{(a^2 + a^2 + a^2)}{(a+a)} = \frac{4}{3\pi} \times \frac{3a^2}{2a} = \frac{2a}{\pi}$$

19 a Use $\bar{x} = \frac{\int y^2 x \, dx}{\int y^2 \, dx}$, and substitute $y = \frac{1}{2x^2}$

$$\begin{aligned}\bar{x} &= \frac{\int_1^2 xy^2 \, dx}{\int_1^2 y^2 \, dx} = \frac{\int_1^2 x \times \frac{1}{4} x^{-4} \, dx}{\int_1^2 \frac{1}{4} x^{-4} \, dx} = \frac{\int_1^2 x^{-3} \, dx}{\int_1^2 x^{-4} \, dx} \\ &= \frac{\left[-\frac{1}{2} x^{-2} \right]_1^2}{\left[-\frac{1}{3} x^{-3} \right]_1^2} = \frac{3(1^{-2} - 2^{-2})}{2(1^{-3} - 2^{-3})} = \frac{3}{2} \times \frac{3}{4} \times \frac{8}{7} = \frac{9}{7}\end{aligned}$$

The centre of mass is $\frac{9}{7}$ m from the y -axis, hence $\left(\frac{9}{7} - 1\right) = \frac{2}{7}$ m from the larger plane face.

b Let ρ be the mass per unit volume of the hemisphere H and the solid S .

$$\text{Then the mass of the } S = \int_1^2 \pi y^2 \rho \, dx = \int_1^2 \frac{\pi}{4} x^{-4} \rho \, dx = \left[-\frac{\pi \rho}{12} x^{-3} \right]_1^2 = \frac{\pi \rho}{12} \left(1 - \frac{1}{8} \right) = \frac{7\pi \rho}{96}$$

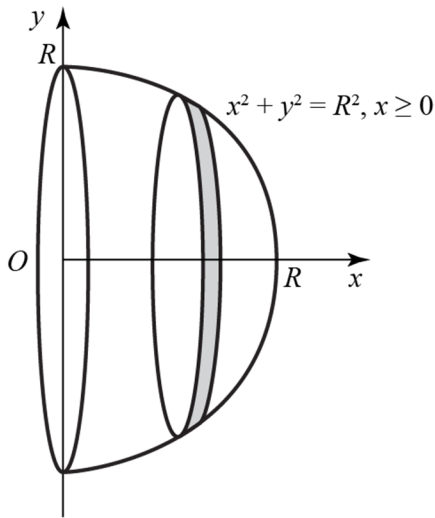
The mass of the hemisphere and its distance of its centre of mass from its plane face can be found from the standard formulae, using $r = 0.5$.

Shape	Mass	Distance of centre of mass from trophy's plane face
Solid, S	$\frac{7\pi\rho}{96}$	$1 - \frac{2}{7} = \frac{5}{7}$
Hemisphere, H	$\frac{2\pi\rho}{3} \left(\frac{1}{2}\right)^3 = \frac{\pi\rho}{12}$	$1 + \frac{3}{8} \times \frac{1}{2} = \frac{19}{16}$
Trophy, T	$\pi\rho \left(\frac{7}{96} + \frac{1}{12}\right) = \frac{5\pi\rho}{32}$	\bar{x}

Taking moments around the trophy's plane face (the bottom of the trophy as shown):

$$\begin{aligned}\frac{5\pi\rho}{32} \bar{x} &= \frac{\pi\rho}{12} \times \frac{19}{16} + \frac{7\pi\rho}{96} \times \frac{5}{7} \\ \bar{x} &= \frac{32}{5} \left(\frac{19}{192} + \frac{5}{96} \right) = \frac{32}{5} \times \frac{29}{192} = \frac{29}{30} = 0.967 \text{ m (3 s.f.)}\end{aligned}$$

20 a A hemisphere is generated when a semicircle is rotated through 180° about the x -axis.



Divide the hemisphere into circular discs, with each disc having mass $\rho\pi y^2 \delta x$ and centre of mass at a distance x from O .

$$\begin{aligned} \text{So } \bar{x} &= \frac{\int_0^R \rho\pi x y^2 dx}{\int_0^R \rho\pi y^2 dx} = \frac{\int_0^R x(R^2 - x^2) dx}{\int_0^R (R^2 - x^2) dx} \\ &= \frac{\left[\frac{1}{2} R^2 x^2 - \frac{1}{4} x^4 \right]_0^R}{\left[R^2 x - \frac{1}{3} x^3 \right]_0^R} = \frac{\frac{1}{2} R^4 - \frac{1}{4} R^4}{R^3 - \frac{1}{3} R^3} = \frac{3}{2} \times \frac{1}{4} R = \frac{3}{8} R \end{aligned}$$

b Using the standard formulae for a cone:

Shape	Mass	Distance of centre of mass from V
Cone	$\frac{1}{3} \pi a^2 k a \rho$	$\frac{3}{4} k a$
Hemisphere	$\frac{2}{3} \pi a^3 \rho$	$k a + \frac{3}{8} a$
Top	$\frac{1}{3} \pi a^3 \rho(k+2)$	\bar{x}

Taking moments about V :

$$\begin{aligned} \frac{1}{3} \pi a^3 \rho(k+2)\bar{x} &= \frac{1}{3} \pi a^3 \rho k \left(\frac{3}{4} k a \right) + \frac{2}{3} \pi a^3 \rho \left(k a + \frac{3a}{8} \right) \\ (k+2)\bar{x} &= \frac{3}{4} k^2 a + 2k a + \frac{3a}{4} \\ \bar{x} &= \frac{(3k^2 + 8k + 3)a}{4(k+2)} \end{aligned}$$

20 c The manufacturer's requirement is that $\bar{x} = ka$

$$\text{Hence from part b } \frac{3k^2 + 8k + 3}{4(k+2)} = k$$

$$\Rightarrow 3k^2 + 8k + 3 = 4k^2 + 8k$$

$$\Rightarrow k^2 = 3, \text{ so } k = \sqrt{3}$$

21 a Using the standard results for a hemisphere:

Shape	Mass	Ratio of masses	Distance of centre of mass from O
Large hemisphere	$\frac{2}{3}\pi a^3 \rho$	8	$\frac{3}{8}a$
Small hemisphere	$\frac{2}{3}\pi \left(\frac{a}{2}\right)^3 \rho$	1	$\frac{3}{16}a$
Remainder	$\frac{2}{3}\pi \frac{7a^3}{8} \rho$	7	\bar{x}

Taking moments about O :

$$7\bar{x} = 8 \times \frac{3}{8}a - 1 \times \frac{3}{16}a$$

$$\bar{x} = \frac{1}{7} \times \frac{45}{16}a = \frac{45a}{112}$$

b Using the result from part a:

	Mass ratios	Distance of centre of mass from O
Bowl	M	$\frac{45}{112}a$
Liquid	kM	$\frac{3}{16}a$
Bowl + liquid	$(k+1)M$	$\frac{17}{48}a$

Taking moments about O :

$$(k+1)M \times \frac{17}{48}a = M \times \frac{45}{112}a + kM \times \frac{3}{16}a$$

$$k \left(\frac{17}{48} - \frac{3}{16} \right) = \frac{45}{112} - \frac{17}{48}$$

$$\frac{8}{48}k = \frac{45}{112} - \frac{17}{48}$$

$$k = 6 \times \left(\frac{45}{112} - \frac{17}{48} \right) = \frac{6(2160 - 1906)}{5376} = \frac{254}{896} = \frac{2}{7}$$

$$22 \text{ a } V = \int_0^1 \pi y^2 dx = \int_0^1 \frac{\pi}{4} (x-2)^4 dx \quad \text{substituting } y = \frac{1}{2}(x-2)^2$$

$$\Rightarrow V = \left[\frac{\pi}{20} (x-2)^5 \right]_0^1 = -\frac{\pi}{20} (-2)^5 = \frac{32\pi}{20} = \frac{8\pi}{5} \text{ cm}^3$$

b Let \bar{x} be the distance of the centre of mass of S from its plane face

To find, using $M\bar{x} = \int \rho\pi y^2 x dx$, and substituting $y^2 = \frac{1}{4}(x-2)^2$ and $M = \frac{8\pi\rho}{5}$ from part a gives:

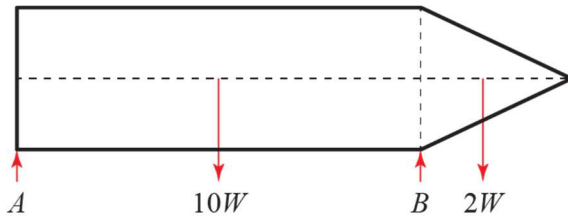
$$\bar{x} = \frac{5}{32} \int_0^2 (x-2)^4 x dx$$

Integrate using the substitution $u = x - 2$

$$\begin{aligned} \bar{x} &= \frac{5}{32} \int_{-2}^0 u^4 (u+2) du = \frac{5}{32} \int_{-2}^0 u^5 + 2u^4 du \\ &= \frac{5}{32} \left[\frac{1}{6} u^6 + \frac{2}{5} u^5 \right]_{-2}^0 = \frac{5}{32} \left(\frac{2}{5} \times 2^5 - \frac{1}{6} 2^6 \right) = \frac{5}{32} \left(\frac{64}{5} - \frac{64}{6} \right) \\ &= 10 \left(\frac{1}{5} - \frac{1}{6} \right) = \frac{10}{30} = \frac{1}{3} \end{aligned}$$

The centre of mass lies on the axis of symmetry at a distance of $\frac{1}{3}$ cm from the plane base.

c Let the reaction force at A be A , and at B be B .



Taking moments about B :

$$A \times 8 + 2W \times \frac{1}{3} = 10W \times 4$$

$$A = \frac{\left(40 - \frac{2}{3}\right)W}{8} = \frac{118W}{24} = \frac{59W}{12}$$

23 a Using the formulae for standard uniform bodies:

Shape	Mass	Mass ratios	Distance of centre of mass from base
Cylinder	$\pi r^2 h \rho$	1	$\frac{h}{2}$
Cone	$\frac{1}{3} \pi r^2 \frac{h}{2} \rho$	$\frac{1}{6}$	$h - \frac{1}{4} \left(\frac{h}{2} \right) = \frac{7h}{8}$
Ornament	$\frac{5}{6} \pi r^2 h \rho$	$\frac{5}{6}$	\bar{x}

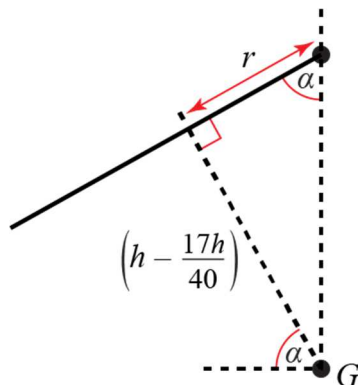
Take moments about O , the centre of plane base:

$$\frac{5}{6} \bar{x} = 1 \times \frac{h}{2} - \frac{1}{6} \times \frac{7h}{8}$$

$$\frac{5}{6} \bar{x} = \frac{h}{2} - \frac{7h}{48} = \frac{17h}{48}$$

$$\bar{x} = \frac{17h}{48} \times \frac{6}{5} = \frac{17h}{40}$$

b The centre of mass G of the ornament will be directly below the point of suspension.



$$\tan \alpha = \frac{h - \frac{17h}{40}}{r} = \frac{23h}{40r}$$

As $h = 4r$, this gives $\tan \alpha = \frac{23r}{10r} = 2.3$

$$\Rightarrow \alpha = 66.5^\circ \text{ (1 d.p.)}$$

- 24 a** Let ρ be the mass per unit area of the material and \bar{x} the distance of the centre of mass of the closed container from O . Using the formulae for standard uniform bodies:

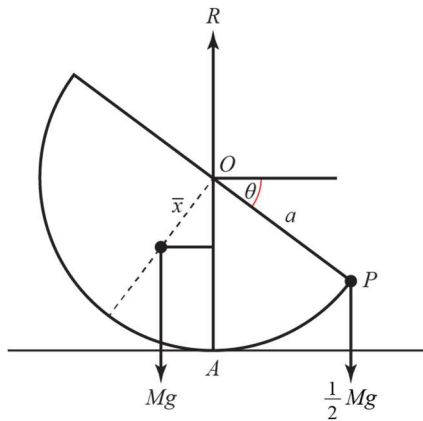
Shape	Mass	Distance of centre of mass from O
Circular disc	$\pi a^2 \rho$	0
Hemispherical bowl	$2\pi a^2 \rho$	$\frac{1}{2}a$
Closed container	$3\pi a^2 \rho$	\bar{x}

Taking moments about O :

$$3\pi a^2 \rho \bar{x} = 0 + 2\pi a^2 \rho \times \frac{a}{2}$$

$$\bar{x} = \frac{a}{3}$$

- b** The container is resting in equilibrium, so the weight of P acts to one side and the weight of C balances on the other side.



Taking moments about O :

$$Mg \times \frac{a}{3} \sin \theta = \frac{1}{2} Mg \times a \cos \theta$$

$$\text{So } \tan \theta = \frac{3}{2}$$

$$\Rightarrow \theta = 56^\circ \text{ (to the nearest degree)}$$

- 25 a** The hemisphere K has mass M . The hemisphere H has double the radius, so it has 8 times the volume. Its mass is $8M$. Therefore the composite body S has mass $M + 8M = 9M$.

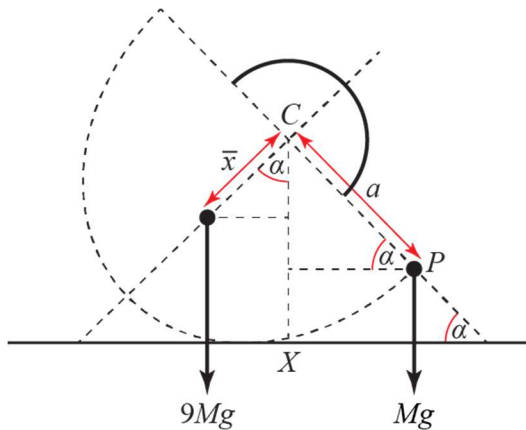
Shape	Mass	Distance of centre of mass from C
H	$8M$	$\frac{3a}{8}$
K	M	$\frac{-3a}{16}$
S	$M + 8M$	\bar{x}

Taking moments about C :

$$9M\bar{x} = 8M \times \frac{3a}{8} - M \times \frac{3a}{16} = M \times \frac{45a}{16}$$

$$\bar{x} = \frac{45a}{9 \times 16} = \frac{5a}{16}$$

- b** The composite body with particle P attached rests in equilibrium, with C above the point of contact with the plane, X .



$$Mg \times a \cos \alpha = 9Mg\bar{x} \sin \alpha = 9 \times \frac{5}{16} Mg \sin \alpha$$

substituting result for \bar{x} from part **a**

$$\Rightarrow \tan \alpha = \frac{16}{45}$$

- 26 a** Let ρ be the mass per unit volume of the material of the cylinder and \bar{x} the distance of the centre of mass of the toy from O . Using the formulae for standard uniform bodies:

Shape	Mass	Mass ratio	Distance of centre of mass from O
Hemisphere	$\frac{2}{3}\pi r^3 \rho$	$4r$	$\frac{5r}{8}$
Cylinder	$\pi r^2 h \rho$	h	$\frac{h}{2} + r$
Toy	$\pi r^2 \rho(4r + h)$	$4r + h$	\bar{x}

Taking moments about O :

$$(4r + h)\bar{x} = 4r \times \frac{5r}{8} + h \left(\frac{h}{2} + r \right)$$

$$\bar{x} = \frac{5r^2 + h^2 + 2rh}{2(4r + h)} = \frac{h^2 + 2hr + 5r^2}{2(h + 4r)} \quad \text{as required}$$

- b** If the toy remains in equilibrium when resting on any point of the curved surface, its centre of mass must be in the centre of the hemisphere's flat surface, so $\bar{x} = r$.

Hence from part **a** $\frac{h^2 + 2hr + 5r^2}{2(h + 4r)} = r$

$$\Rightarrow h^2 + 2hr + 5r^2 = 2rh + 8r^2$$

$$\Rightarrow h^2 = 3r^2$$

$$\Rightarrow h = \sqrt{3}r$$

- 27 a** Let \bar{x} be distance of the centre of mass from AB on the axis of symmetry in the direction away from O . Using the formulae for standard uniform bodies:

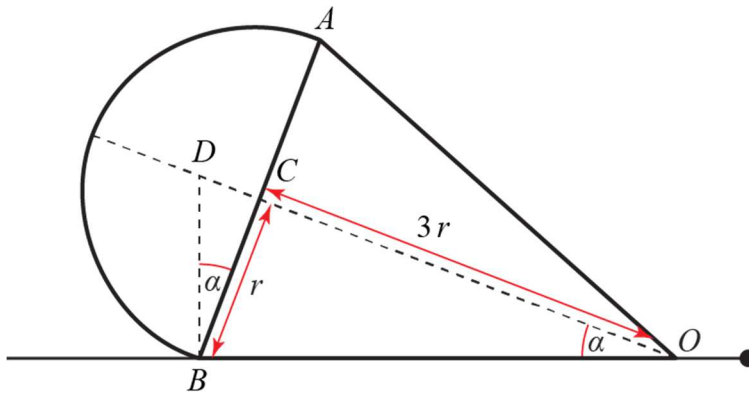
Shape	Mass	Distance of centre of mass from AB
Hemisphere	M	$\frac{3}{8}r$
Cone	m	$-\frac{3}{4}r$
Toy	$m + M$	\bar{x}

Taking moments about AB :

$$(m + M)\bar{x} = \frac{3}{8}Mr - \frac{3}{4}mr = \frac{3r}{8}(M - 2m)$$

$$\bar{x} = \frac{3(M - 2m)}{8(M + m)}r \quad \text{as required}$$

- 27 b** Let D be the point on the axis of symmetry vertically above B when the toy is placed on a horizontal surface. The toy will not remain in equilibrium if its centre of mass is not on the line segment OD , i.e if $\bar{x} > CD$.



$$\text{From the diagram: } \tan \alpha = \frac{r}{3r} = \frac{CD}{r} \Rightarrow CD = \frac{1}{3}r$$

$$\text{So } \bar{x} > CD \Rightarrow \frac{3(M-2m)}{8(M+m)}r > \frac{1}{3}r$$

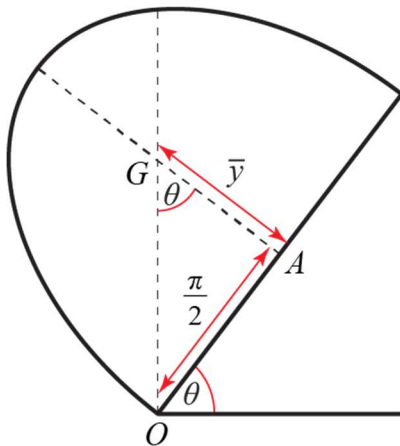
$$\Rightarrow 9(M-2m) > 8(M+m)$$

$$\Rightarrow M > 26m \quad \text{as required}$$

- 28 a** Using $\bar{y} = \frac{\int \frac{1}{2}y^2 dx}{\int y dx}$ with $y = \sin x$ gives:

$$\begin{aligned} \bar{y} &= \frac{\int_0^\pi \frac{1}{2} \sin^2 x dx}{\int_0^\pi \sin x dx} = \frac{\frac{1}{4} \int_0^\pi (1 - \cos 2x) dx}{\int_0^\pi \sin x dx} \\ &= \frac{1}{4} \frac{[x - \frac{1}{2} \sin 2x]_0^\pi}{[-\cos x]_0^\pi} = \frac{1}{4} \frac{\pi}{(1+1)} = \frac{\pi}{8} \end{aligned}$$

- b** When S is on the point of toppling, G is above O . Let A be the point midway along the base of S .



$$\tan \theta = \frac{\frac{\pi}{2}}{\frac{\pi}{8}} = \frac{\pi}{2} = 4$$

$$\Rightarrow \theta = 76^\circ \quad (\text{to the nearest degree})$$

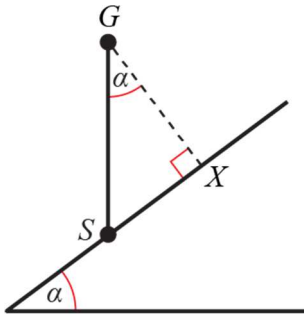
- 29 a** Let ρ be the mass per unit volume of the material and \bar{x} be the distance of the centre of mass of the solid S from O . Using the formulae for standard uniform bodies:

Shape	Mass	Mass ratios	Distance of centre of mass from O
Cylinder	$\pi\rho(2a)^2\left(\frac{3}{2}a\right)$	6	$\frac{3}{4}a$
Hemisphere	$\frac{2}{3}\pi\rho a^3$	$\frac{2}{3}$	$\frac{3}{8}a$
Solid, S	$\pi\rho\left(6a^3 - \frac{2}{3}a^3\right)$	$\frac{16}{3}$	\bar{x}

Taking moments about O :

$$\begin{aligned}\frac{16}{3}\bar{x} &= 6 \times \frac{3}{4}a - \frac{2}{3} \times \frac{3}{8}a \\ &= \frac{9}{2}a - \frac{1}{4}a = \frac{17}{4}a \\ \bar{x} &= \frac{51a}{64} = 0.797a \quad (3 \text{ s.f.})\end{aligned}$$

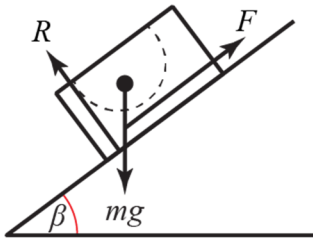
- b** On the point of toppling: G is above S – the lowest point on the bottom circular face.



Let X be the centre of the base of the cylinder.

$$\begin{aligned}\tan \alpha &= \frac{SX}{XG} = \frac{2a}{\frac{3}{2}a - \bar{x}} = \frac{64 \times 2}{96 - 51} = \frac{128}{45} \quad \text{substituting value for } \bar{x} \text{ from part a} \\ \Rightarrow \alpha &= 70.6^\circ \quad (3 \text{ s.f.})\end{aligned}$$

29 c There are three forces acting on S , its weight, the reaction force and friction.



When S is on the point of sliding $F = \mu R = 0.8R$

Resolving perpendicular to the plane

$$R (\perp): R - mg \cos \beta = 0 \Rightarrow R = mg \cos \beta$$

Resolving along the plane

$$R (\parallel): F - mg \sin \beta = 0 \Rightarrow F = mg \sin \beta$$

Using $F = 0.8R$ gives

$$mg \sin \beta = 0.8mg \cos \beta \Rightarrow \tan \beta = 0.8$$

So $\beta = 38.7^\circ$ (3 s.f.)

30 a Using the formulae for standard uniform bodies:

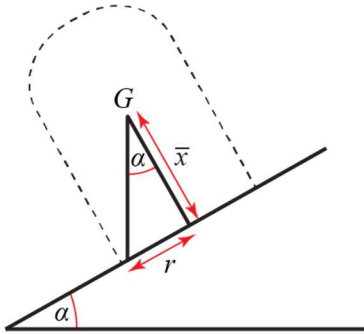
Shape	Mass	Distance of centre of mass from O
Hemisphere, H	$2M$	$h + \frac{3}{8}r$
Cylinder, C	$3M$	$\frac{h}{2}$
Body	$5M$	\bar{x}

Taking moments about O :

$$5M\bar{x} = 2M\left(h + \frac{3}{8}r\right) + 3M \times \frac{h}{2} = 2h + \frac{3}{4}r + \frac{3h}{2} = \frac{7h}{2} + \frac{3r}{4}$$

$$\bar{x} = \frac{14h + 3r}{20}$$

- 30 b** When the body is on the point of toppling, its centre of mass G lies directly above a point on the plane which is at the lowest point on its cylindrical base.



$$\text{From the diagram: } \tan \alpha = \frac{r}{\bar{x}} = \frac{20r}{14h + 3r}$$

$$\text{As } \tan \alpha = \frac{4}{3}, \text{ this gives } \frac{20r}{14h + 3r} = \frac{4}{3}$$

$$\Rightarrow 60r = 56h + 12r$$

$$\Rightarrow 48r = 56h$$

$$\Rightarrow h = \frac{48}{56}r = \frac{6}{7}r$$

- 31 a** Let ρ be the mass per unit volume of the material in the cylinder and \bar{x} be the distance of the centre of mass of the top from O . Using the formulae for standard uniform bodies:

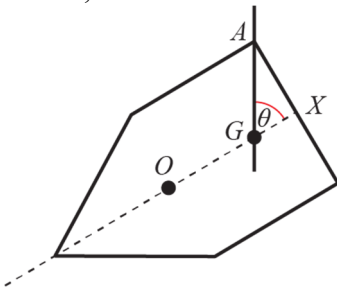
Shape	Mass	Mass ratio	Position of centre of mass from O
Cylinder	$36\pi\rho r^3$	1	$2r$
Cone	$36\pi\rho r^3$	1	$-r$
Toy	$72\pi\rho r^3$	2	\bar{x}

Taking moments about O :

$$2\bar{x} = 1 \times 2r + 1 \times (-r) = r$$

$$\bar{x} = \frac{r}{2}$$

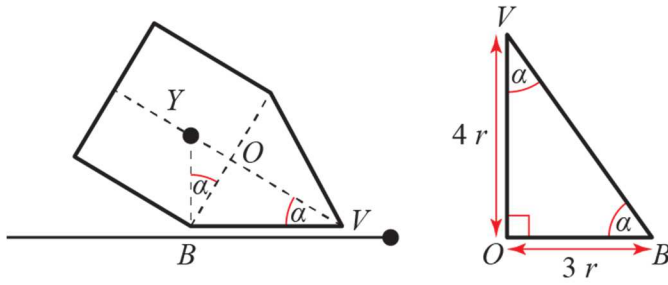
- b** In $\triangle AGX$, AX is the radius $3r$ of the cylinder and X is the centre of the base of the cylinder.



$$\tan \theta = \frac{AX}{GX} = \frac{AX}{OX - OG} = \frac{3r}{4r - \bar{x}} = \frac{3r}{4r - \frac{r}{2}} = \frac{6}{7}$$

$$\Rightarrow \theta = 47.5^\circ \text{ (1 d.p.)}$$

- 31 c The toy will not topple if G (the centre of mass) is vertically above a point between B and V . Let Y be the point directly above B on the axis of symmetry of the toy.



$$\tan \alpha = \frac{3r}{4r} \text{ from } \triangle VOB \text{ and } \tan \alpha = \frac{OY}{OB} = \frac{OY}{3r} \text{ from } \triangle BOY$$

$$\text{So } \frac{OY}{3r} = \frac{3}{4} \Rightarrow OY = \frac{9r}{4}$$

$$\text{But } OG = \frac{r}{2} \text{ so } OG < OY$$

This means that the centre of mass is above the face of the body in contact with the place, so the toy will not topple.

Challenge

- 1 The two particles have equal but opposite velocities just before collision at point A . Let their velocities at point A be u and $-u$ respectively.

After collision, P_1 travels three times the distance travelled by P_2 before colliding at point B . Therefore, the velocities after collision are $-3v$ and v respectively.

Using conservation of momentum:

$$mu + 2m(-u) = m(-3v) + 2mv$$

$$\Rightarrow u = v$$

$$\text{Coefficient of restitution: } e = \frac{u - (-u)}{u - (-3u)} = \frac{1}{2}$$

Challenge

- 2 a Let the instantaneous speed at any point be w and the normal reaction between the ball and the ring is N . Then resolving towards the centre of the circle:

$$N = \frac{mw^2}{R} \quad \text{using Newton's second law}$$

So the frictional force opposite to the direction of motion is given by:

$$F = \mu \frac{mw^2}{R}$$

Resolving in the direction of motion:

$$m \frac{dw}{dt} = -\mu \frac{mw^2}{R}$$

Let v be the velocity at time t , so integrating:

$$\int_u^v \frac{1}{w^2} dw = -\frac{\mu}{R} \int_0^t dt$$

$$\Rightarrow \left[-\frac{1}{w} \right]_u^v = -\frac{\mu}{R} [t]_0^t$$

$$\Rightarrow -\frac{1}{v} + \frac{1}{u} = -\frac{\mu}{R} t$$

$$\Rightarrow \frac{1}{v} = \frac{1}{u} + \frac{\mu}{R} t = \frac{R + u\mu t}{uR}$$

$$\Rightarrow v = \frac{uR}{R + u\mu t}$$

- b From part a $v = \frac{dx}{dt} = \frac{uR}{R + u\mu t}$

The ball completes one revolution in time t when $x = 2\pi R = \pi$

$$\int_0^\pi 1 dx = \int_0^t \frac{uR}{R + u\mu t} dt \Rightarrow \int_0^t \frac{20}{0.5 + 10t} dt$$

$$\Rightarrow \pi = [2 \ln(0.5 + 10t)]_0^t = 2(\ln(0.5 + 10t) - \ln 0.5)$$

$$\Rightarrow \pi = 2 \ln(1 + 20t)$$

$$\Rightarrow 20t = e^{\frac{\pi}{2}} - 1$$

$$\Rightarrow t = \frac{e^{\frac{\pi}{2}} - 1}{20} = 0.191 \text{ s (3 s.f.)}$$